

Sinos de Água

Prof. Cassio Sigaud

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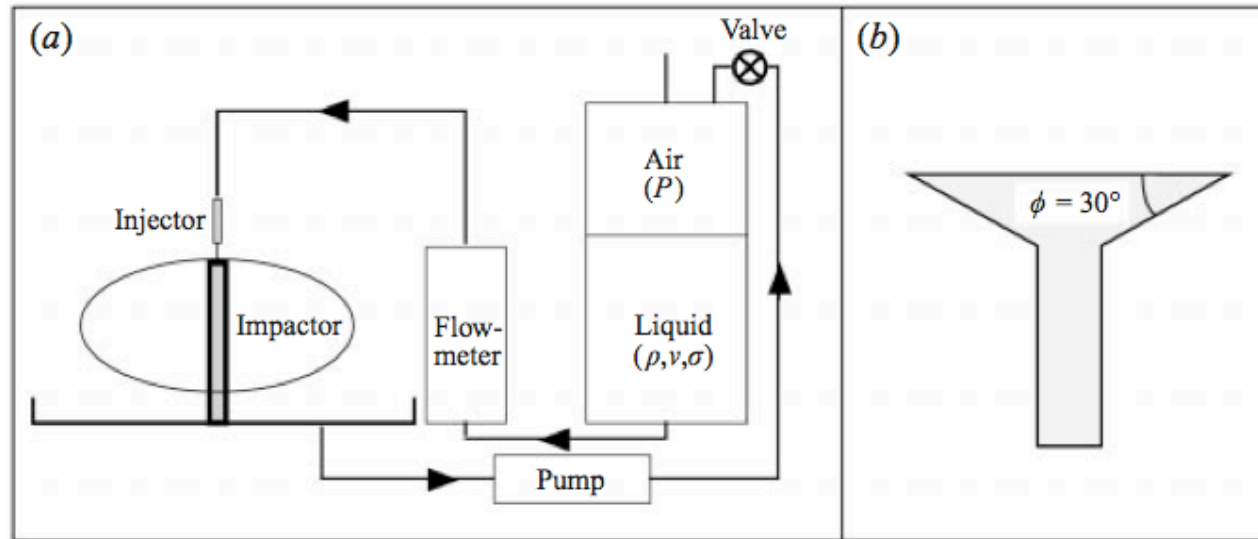


FIGURE 3. Presentation of the experimental set-up: (a) general view, (b) detail of the impactor.

Dynamics and stability of water bells

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Stability of Water Bells Generated by Jet Impacts on a Disk

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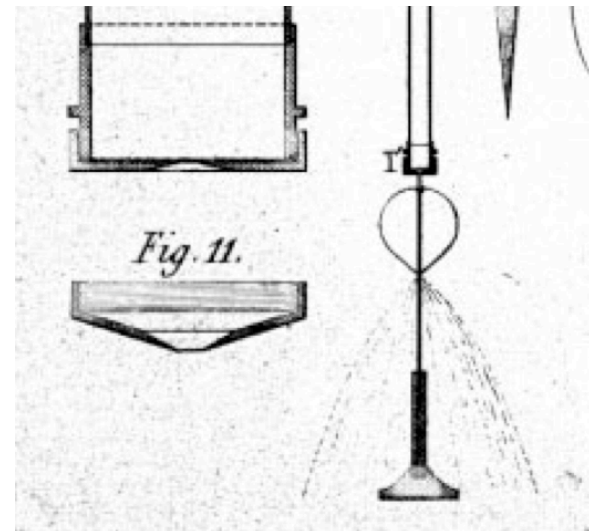
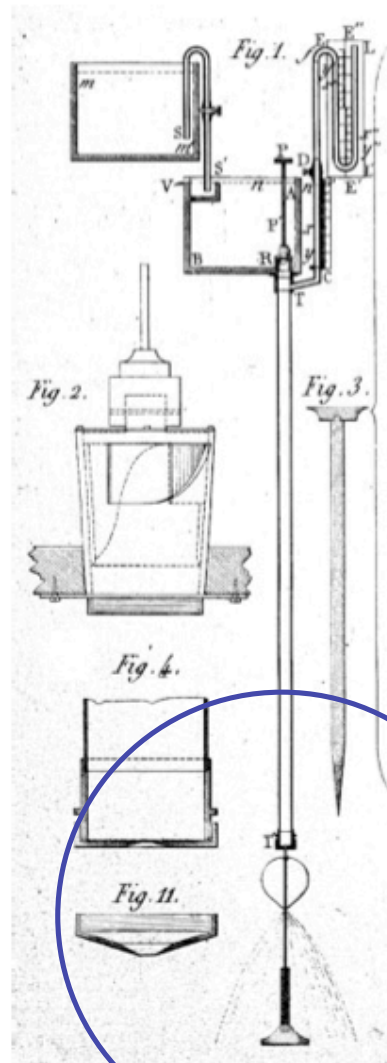


FIGURE 4. Experimental setup used by Savart (from Planche 4 of Savart's 1833 papers).

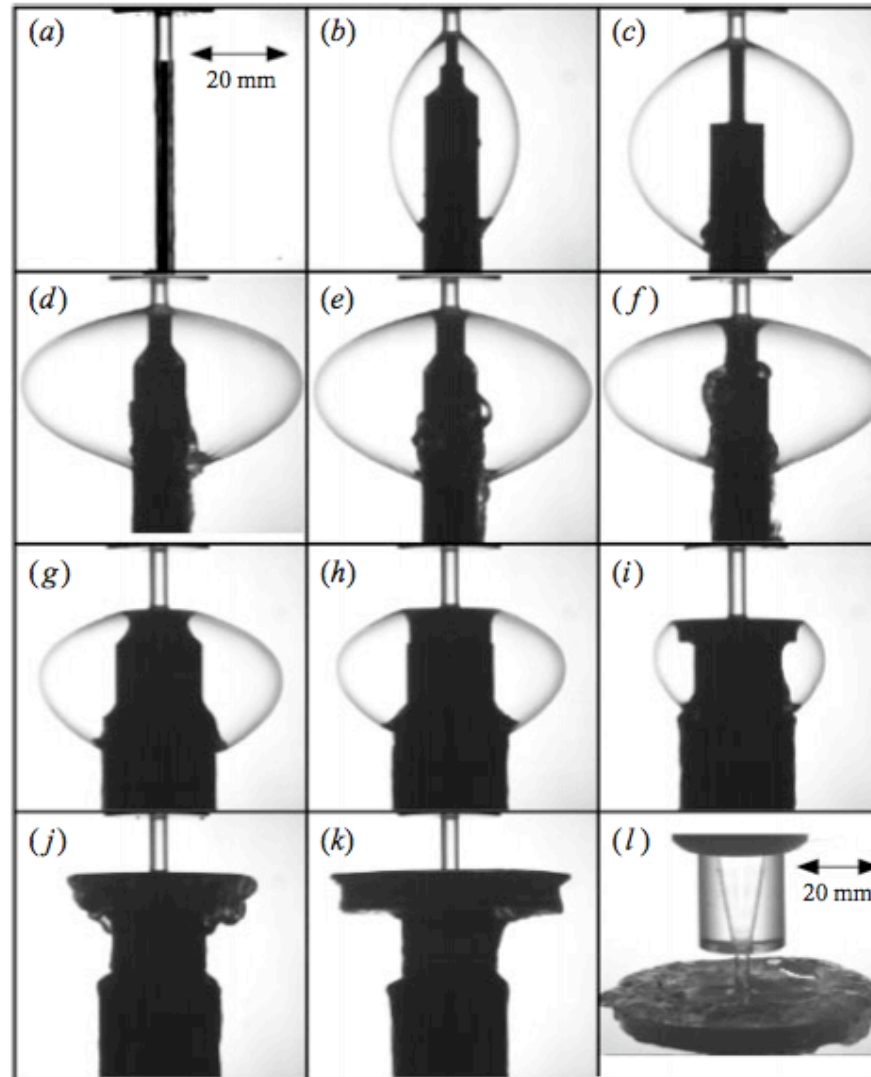


FIGURE 5. Influence of the impact diameter, D_i , for $D_0 = 3$ mm, $U_0 = 2.08$ m s⁻¹, $Re = 6240$, $We = 178$: (a) $D_i = 1.18$ mm, (b) 3.0 mm, (c) 4.0 mm, (d) 5.44 mm, (e) 7.33 mm, (f) 9.87 mm, (g) 13.3 mm, (h) 17.81 mm, (i) 24.13 mm, (j) 32.51 mm, (k) 43.79 mm, (l) 58.99 mm.

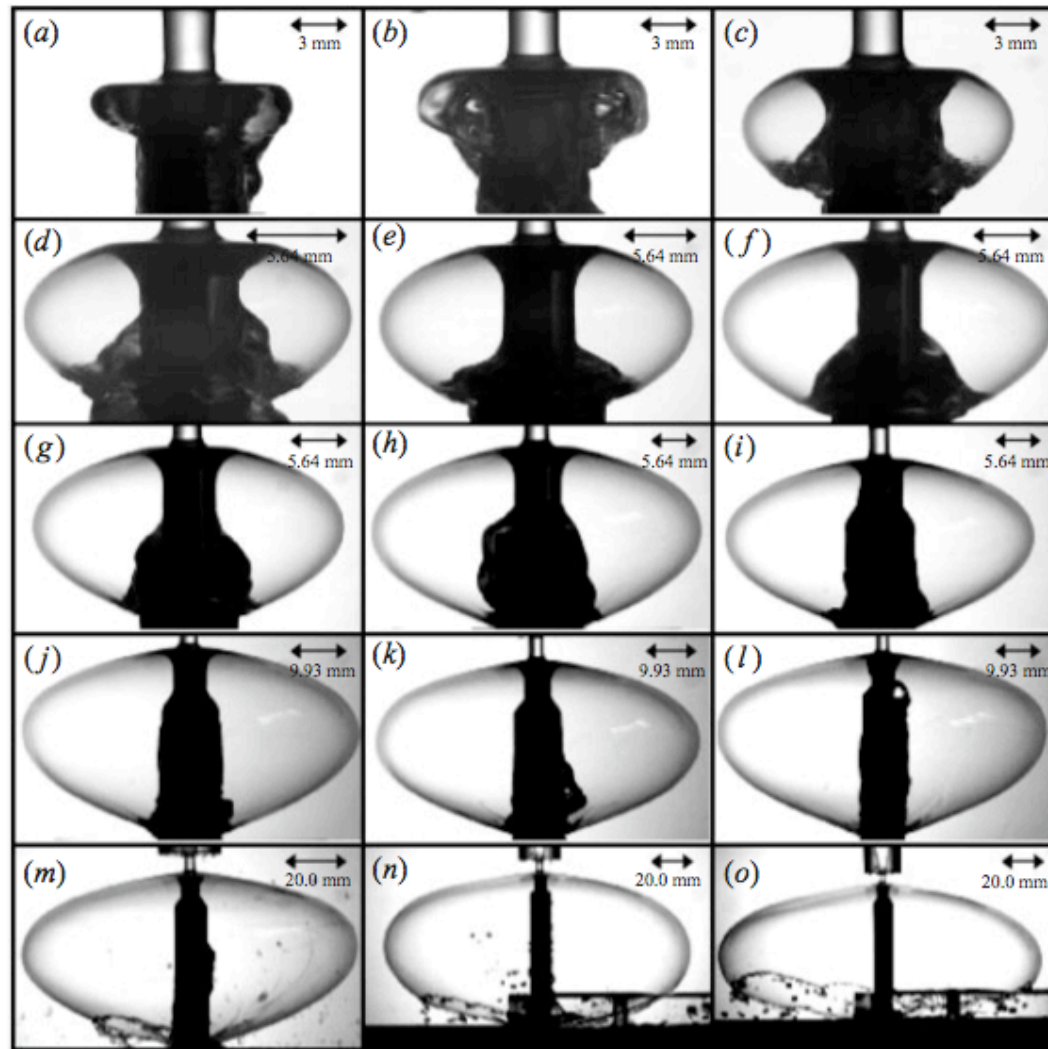


FIGURE 7. Influence of the impact velocity U_0 for $D_0 = 3$ mm, $D_t = 9.87$ mm: (a) $U_0 = 0.63$ m s⁻¹, (b) 1.0 m s⁻¹, (c) 1.08 m s⁻¹, (d) 1.2 m s⁻¹, (e) 1.42 m s⁻¹, (f) 1.57 m s⁻¹, (g) 1.68 m s⁻¹, (h) 1.87 m s⁻¹, (i) 2.08 m s⁻¹, (j) 2.25 m s⁻¹, (k) 2.42 m s⁻¹, (l) 2.65 m s⁻¹, (m) 2.88 m s⁻¹, (n) 3.43 m s⁻¹, (o) 3.67 m s⁻¹.

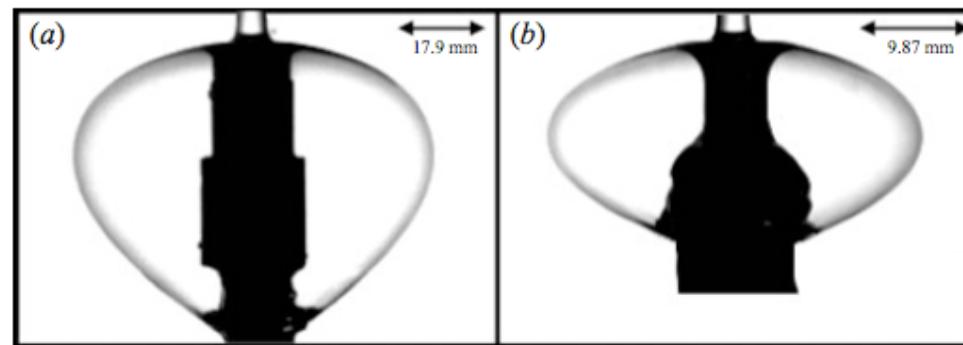


FIGURE 8. Influence of the jet diameter D_0 : (a) $D_0 = 6$ mm, $D_i = 17.91$ mm, $U_0 = 1.19$ m s⁻¹, $\psi_0 = 73.8^\circ$, $X = 2.985$, $We = 116$ and $Re = 7140$; (b) $D_0 = 3$ mm, $D_i = 9.87$ mm, $U_0 = 1.68$ m s⁻¹, $\psi_0 = 72.2^\circ$, $X = 3.29$, $We = 116$ and $Re = 5040$.

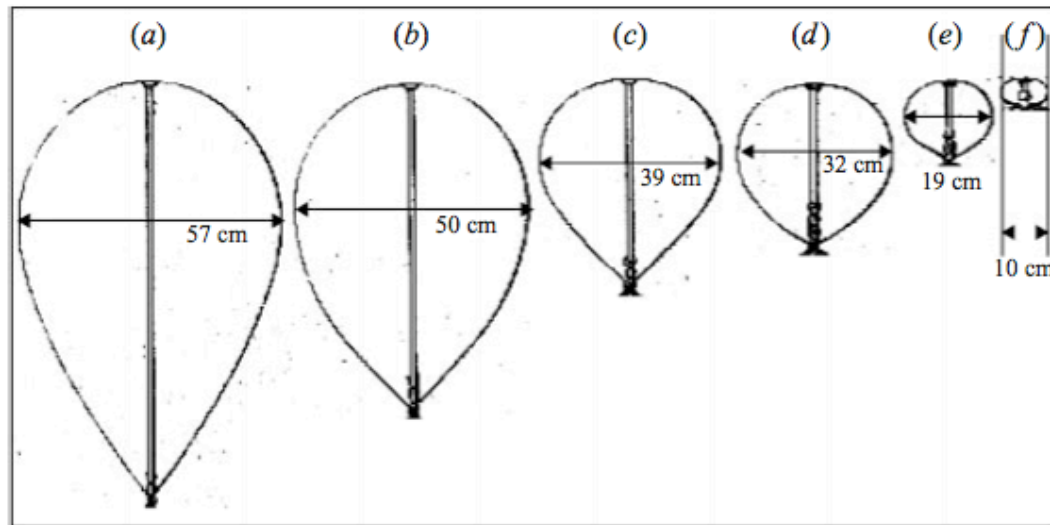


FIGURE 9. Savart's observations of the influence of the jet diameter: (a) $D_0 = 14.4$ mm, $U_0 = 2.06$ m s $^{-1}$, (b) $D_0 = 12$ mm, $U_0 = 2.15$ m s $^{-1}$, (c) $D_0 = 9.6$ mm, $U_0 = 2.54$ m s $^{-1}$, (d) $D_0 = 7.2$ mm, $U_0 = 2.67$ m s $^{-1}$, (e) $D_0 = 4.8$ mm, $U_0 = 3.16$ m s $^{-1}$, (f) $D_0 = 2.4$ mm, $U_0 = 5.66$ m s $^{-1}$.

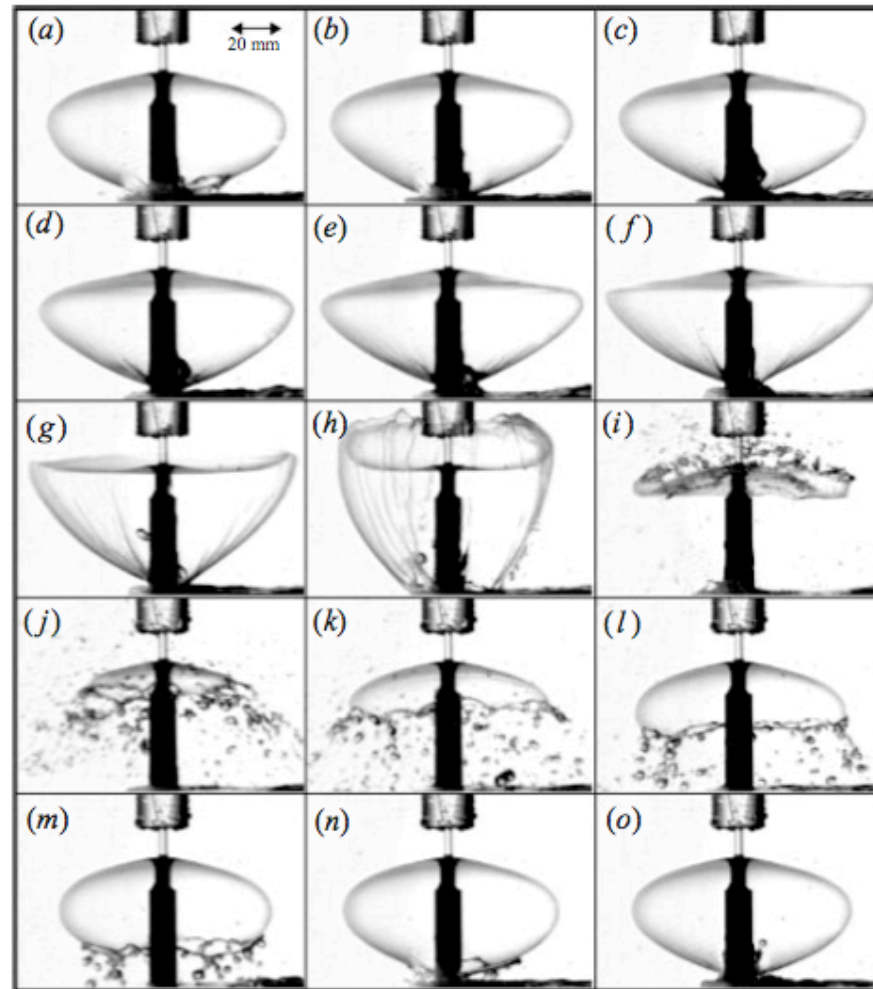


FIGURE 16. Instability of a closed water bell triggered by a pressure perturbation, observed with $\psi_0 = 77^\circ$, $D_0 = 3.0$ mm, $D_i = 9.87$ mm and $U_0 = 2.7$ m s $^{-1}$. Time increases from (a) to (o) with the time step $\Delta t = 35.5$ ms.

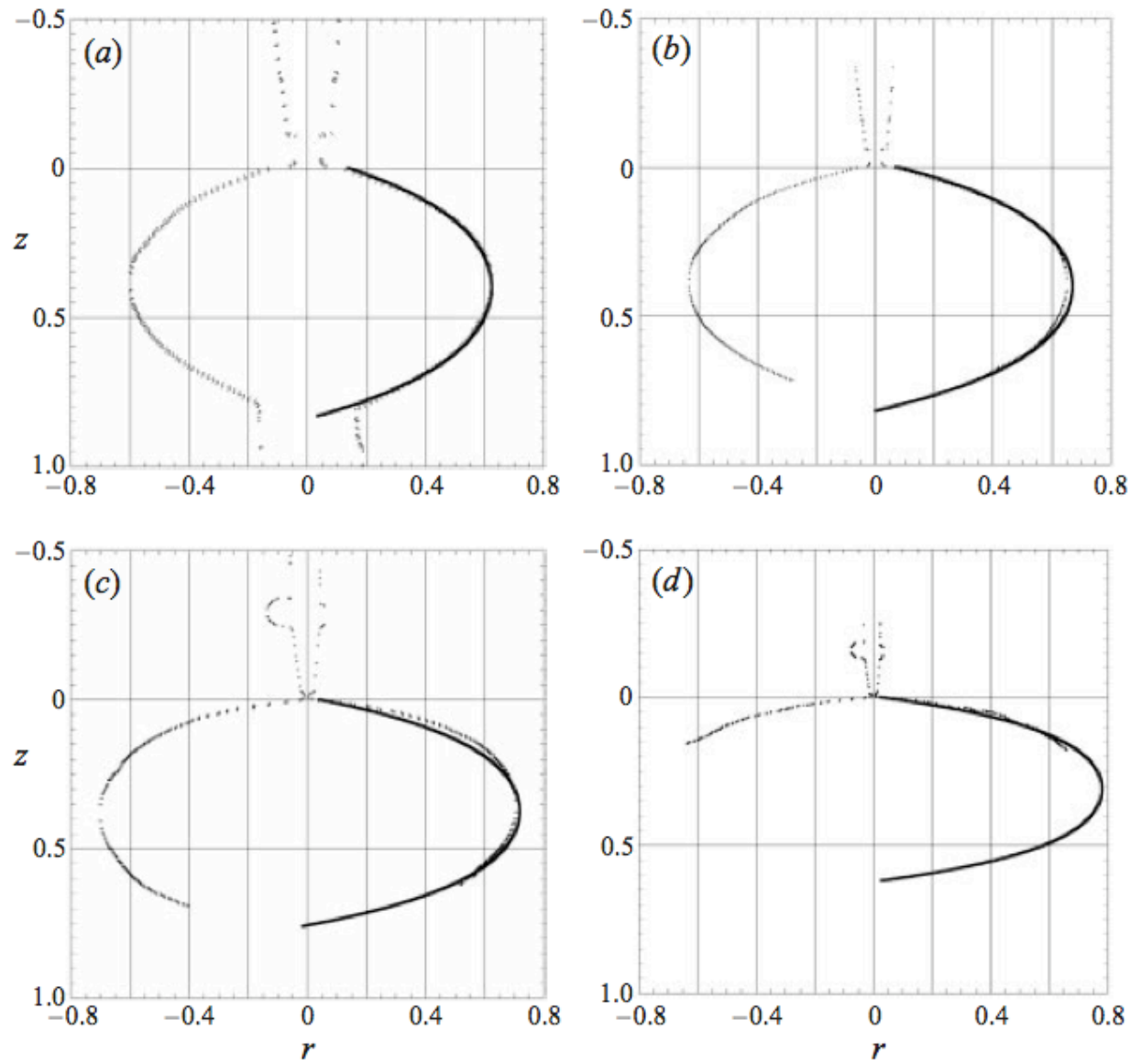


FIGURE 22. Comparison between the measured bell shape and the theoretical catenary: (a) $u_e = 0.84$, $r_i = 0.139$, $\psi_0 = 72^\circ$; (b) $u_e = 0.86$, $r_i = 0.071$, $\psi_0 = 76^\circ$; (c) $u_e = 0.88$, $r_i = 0.040$, $\psi_0 = 79^\circ$; (d) $u_e = 0.89$, $r_i = 0.022$, $\psi_0 = 82.5^\circ$.

terms of X and Y , by integrating $\tan \phi = dY/dX$ using (I 11)

$$X = \cos \phi_0 [\cosh^{-1}(\sec \phi_0) - \cosh^{-1}\{(1 - Y) \sec \phi_0\}]. \quad (\text{I } 12)$$

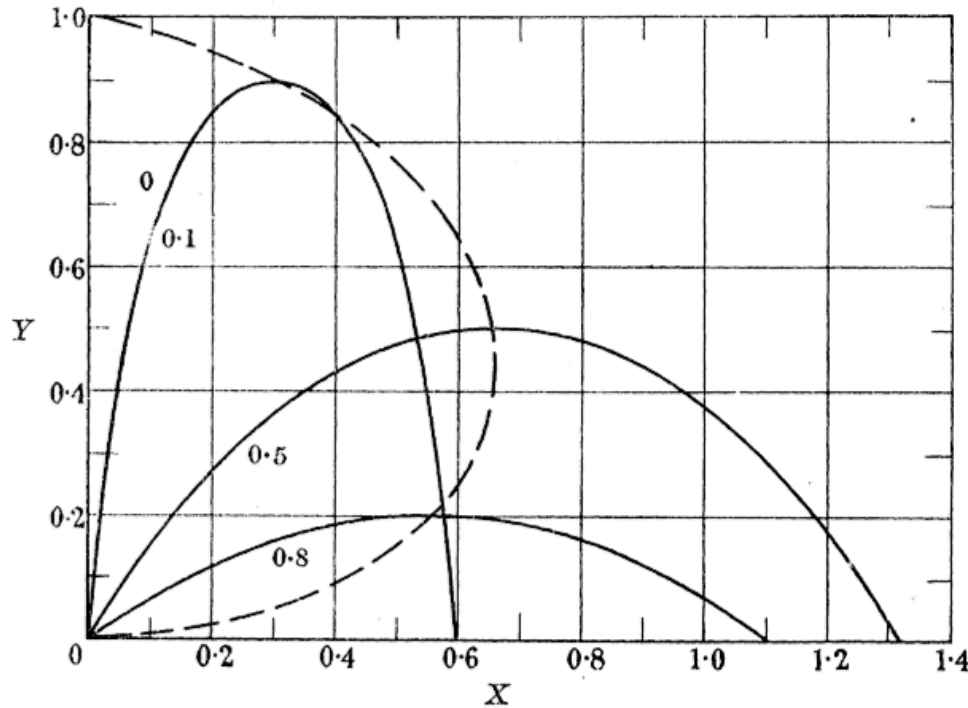


FIGURE I 1. Calculated meridian sections of water bells. The numbers give the values of $\cos \phi_0$.

Expression (I 12) applies in the range $0 < x < \cos \phi_0 \cosh^{-1}(\sec \phi_0)$. At the upper limit $1 - Y = \cos \phi_0$ and $dY/dX = 0$ but d^2Y/dX^2 is finite. The solution can be

The dynamics of thin sheets of fluid
I. Water bells

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